# Polarization Optics 

Nicolas Fressengeas

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## Polarization Optics

## N. Fressengeas

Laboratoire Matériaux Optiques, Photonique et Systèmes Unité de Recherche commune à l'Université de Lorraine et à Supélec

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## Further reading

[Hua94, GB94]

圊 A. Gerrard and J.M. Burch.
Introduction to matrix methods in optics.
Dover, 1994.
囯 S. Huard.
Polarisation de la lumière.
Masson, 1994.

## Course Outline

(1) The physics of polarization optics

- Polarization states
- Jones Calculus
- Stokes parameters and the Poincare Sphere
(2) Polarized light propagation
- Jones Matrices Examples
- Matrix, basis \& eigen polarizations
- Jones Matrices Composition
(3) Partially polarized light
- Formalisms used
- Propagation through optical devices
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三

## The vector nature of light Optical wave can be polarized, sound waves cannot

## The scalar monochromatic plane wave <br> The electric field reads:

- Electric field is orthogonal to wave and Poynting vectors


## The vector nature of light Optical wave can be polarized, sound waves cannot

## The scalar monochromatic plane wave

The electric field reads:

$$
A \cos (\omega t-k z-\varphi)
$$

A vector monochromatic plane wave

- Electric field is orthogonal to wave and Poynting vectors
- Lies in the wave vector normal plane
- Needs 2 components


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A vector monochromatic plane wave

- Electric field is orthogonal to wave and Poynting vectors
- Lies in the wave vector normal plane
- Needs 2 components
- $E_{x}=A_{x} \cos \left(\omega t-k z-\varphi_{x}\right)$
- $E_{y}=A_{y} \cos \left(\omega t-k z-\varphi_{y}\right)$

Polarization states

## Linear and circular polarization states

In phase components $\quad \varphi_{y}=\varphi_{x}$


## Linear and circular polarization states

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Polarization states

## Linear and circular polarization states

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Polarization Optics, version 2.0, frame 5

Polarization states
Jones Calculus
Stokes parameters and the Poincare Sphere

## Linear and circular polarization states

## In phase components <br> $\varphi_{y}=\varphi_{x}$ <br> 

$\pi / 2$ shift $\quad \varphi_{y}=\varphi_{x} \pm \pi / 2$

$$
\pi \text { shift } \quad \varphi_{y}=\varphi_{x}+\pi
$$


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## Left or Right

Polarization Optics, version 2.0, frame 5

Polarization states

## Linear and circular polarization states

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$\pi / 2$ shift $\quad \varphi_{y}=\varphi_{x} \pm \pi / 2$


Left or Right

Polarization Optics, version 2.0, frame 5

Polarization states
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## Linear and circular polarization states

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$\pi / 2$ shift
$\varphi_{y}=\varphi_{x} \pm \pi / 2$


## Left or Right

Polarization Optics, version 2.0, frame 5

## The elliptic polarization state

The polarization state of ANY monochromatic wave

$$
\varphi_{y}-\varphi_{x}= \pm 0
$$



## Electric field

- $E_{x}=A_{x} \cos \left(\omega t-k z-\varphi_{x}\right)$
- $E_{y}=A_{y} \cos \left(\omega t-k z-\varphi_{y}\right)$


## The elliptic polarization state

The polarization state of ANY monochromatic wave

$$
\varphi_{y}-\varphi_{x}= \pm \pi / 8
$$



## Electric field

- $E_{x}=A_{x} \cos \left(\omega t-k z-\varphi_{x}\right)$
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## Electric field

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## The elliptic polarization state

## The polarization state of ANY monochromatic wave

$$
\varphi_{y}-\varphi_{x}= \pm 3 \pi / 4
$$



## Electric field

- $E_{x}=A_{x} \cos \left(\omega t-k z-\varphi_{x}\right)$
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## The elliptic polarization state

The polarization state of ANY monochromatic wave

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$$



## Electric field

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## Electric field

- $E_{x}=A_{x} \cos \left(\omega t-k z-\varphi_{x}\right)$
- $E_{y}=A_{y} \cos \left(\omega t-k z-\varphi_{y}\right)$


## 4 real numbers

- $A_{x}, \varphi_{x}$
- $A_{y}, \varphi_{y}$
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Polarization Optics, version 2.0, frame 6

## The elliptic polarization state

The polarization state of ANY monochromatic wave

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## Electric field

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## 4 real numbers

- $A_{x}, \varphi_{x}$
- $A_{y}, \varphi_{y}$


## 2 complex numbers

- $A_{x} \exp \left(-i \varphi_{x}\right)$
- $A_{y} \exp \left(-i \varphi_{y}\right)$

Polarization Optics, version 2.0, frame 6

## The elliptic polarization state

The polarization state of ANY monochromatic wave

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三

## Polarization states are vectors

Monochromatic polarizations belong to a 2D vector space based on the Complex Ring

ANY elliptic polarization state $\Longleftrightarrow$ Two complex numbers
A set of two ordered complex numbers is one 2D complex vector

## Polarization states are vectors

Monochromatic polarizations belong to a 2D vector space based on the Complex Ring

## ANY elliptic polarization state <br> Two complex numbers

A set of two ordered complex numbers is one 2D complex vector

## Canonical Basis

$$
\left(\left[\begin{array}{l}
1 \\
0
\end{array}\right],\left[\begin{array}{l}
0 \\
1
\end{array}\right]\right)
$$


$\square$

```
Reversed circular
```


## Polarization states are vectors

Monochromatic polarizations belong to a 2D vector space based on the Complex Ring

ANY elliptic polarization state $\quad \Longleftrightarrow$ Two complex numbers
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$$

Polarization Basis
Two independent polarizations

## Link with optics ?

- These two vectors represent two polarization states
- We must decide which ones !


## Polarization states are vectors

Monochromatic polarizations belong to a 2D vector space based on the Complex Ring

ANY elliptic polarization state
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## Link with optics ?

- These two vectors represent two polarization states
- We must decide which ones !


## Polarization Basis

Two independent polarizations :

- Crossed Linear
- Reversed circular
- ...
- YOUR choice


## Examples: Linear Polarizations

## Canonical Basis Choice

- $\left[\begin{array}{l}1 \\ 0\end{array}\right]$ : horizontal linear polarization
- $\left[\begin{array}{l}0 \\ 1\end{array}\right]$ : vertical linear polarization


## Examples: Linear Polarizations

## Canonical Basis Choice

- $\left[\begin{array}{l}1 \\ 0\end{array}\right]$ : horizontal linear polarization
- $\left[\begin{array}{l}0 \\ 1\end{array}\right]$ : vertical linear polarization

$$
\begin{array}{lll}
\hline \text { Tilt } & \theta=\pi / 4 \\
& \frac{1}{\sqrt{2}}\left[\begin{array}{l}
1 \\
1
\end{array}\right] &
\end{array}
$$

## Examples: Linear Polarizations

## Canonical Basis Choice

- $\left[\begin{array}{l}1 \\ 0\end{array}\right]$ : horizontal linear polarization
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## Tilt <br> $$
\theta=3 \pi / 4
$$

$$
\frac{1}{\sqrt{2}}\left[\begin{array}{c}
-1 \\
1
\end{array}\right]
$$

## Examples: Linear Polarizations

## Canonical Basis Choice

- $\left[\begin{array}{l}1 \\ 0\end{array}\right]$ : horizontal linear polarization
- $\left[\begin{array}{l}0 \\ 1\end{array}\right]$ : vertical linear polarization


## Tilt

$$
\left[\begin{array}{c}
\cos (\theta) \\
\sin (\theta)
\end{array}\right]
$$



## Examples: Linear Polarizations

## Canonical Basis Choice

- $\left[\begin{array}{l}1 \\ 0\end{array}\right]$ : horizontal linear polarization
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Tilt

$$
\left[\begin{array}{c}
\cos (\theta) \\
\sin (\theta)
\end{array}\right]
$$

Linear polarization Jones vector
in a linear polarization basis
Linear Polarization : two in phase components

## Examples: Circular Polarizations

In the same canonical basis choice : linear polarizations

$$
\varphi_{y}-\varphi_{x}= \pm \pi / 2
$$



## Electric field

- $E_{x}=A_{x} \cos \left(\omega t-k z-\varphi_{x}\right)$
- $E_{y}=A_{y} \cos \left(\omega t-k z-\varphi_{y}\right)$
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Polarization Optics, version 2.0, frame 9

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## Electric field

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## Jones vector

$$
\frac{1}{\sqrt{2}}\left[\begin{array}{c}
1 \\
\pm i
\end{array}\right]
$$

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## About changing basis

A polarization state Jones vector is basis dependent

## Some elementary algebra

- The polarization vector space dimension is 2
- Therefore : two non colinear vectors form a basis
- Any polarization state can be expressed as the sum of two non colinear other states


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## Find the transformation matrix between between the two following

- Horizontal and Vertical Linear Polarizations


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- Any polarization state can be expressed as the sum of two non colinear other states
- Remark : two colinear polarization states are identical


## Homework

Find the transformation matrix between between the two following bases :

- Horizontal and Vertical Linear Polarizations
- Right and Left Circular Polarizations


## Relationship between Jones and Poynting vectors

 Jones vectors also provide information about intensity
## Choose an orthonormal basis

- Hermitian product is null : $\overline{J_{1}} \cdot J_{2}=0$
- Each vector norm is unity : $\overline{J_{1}} \cdot J_{1}=\overline{J_{2}} \cdot J_{2}=1$


## Relationship between Jones and Poynting vectors

 Jones vectors also provide information about intensity
## Choose an orthonormal basis

- Hermitian product is null : $\overline{J_{1}} \cdot J_{2}=0$
- Each vector norm is unity : $\overline{J_{1}} \cdot J_{1}=\overline{J_{2}} \cdot J_{2}=1$

Hermitian Norm is Intensity
Simple calculations show that :

- If each Jones component is one complex electric field component
- The Hermitian norm is proportional to beam intensity
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3 Partially polarized light

- Formalisms used
- Propagation through optical devices


## The Stokes parameters

A set of 4 dependent real parameters that can be measured
$P_{0} \quad$ Overall Intensity
$P_{0}=1$

## The Stokes parameters

A set of 4 dependent real parameters that can be measured
$P_{0}=1$

$$
P_{1}=I_{x}-I_{y}
$$

틀

## The Stokes parameters

A set of 4 dependent real parameters that can be measured

| $P_{0}$ | Overall Intensity |
| :---: | :---: |
|  | $P_{0}=I$ |
| $P_{2}$ | in a $\pi / 4$ Tilted Basis |
|  | $P_{2}=I_{\pi / 4}-I_{-\pi / 4}$ |

## The Stokes parameters

A set of 4 dependent real parameters that can be measured


## Relationship between Jones and Stockes

## Sample Jones Vector

$$
J=\left[\begin{array}{l}
A_{x} \exp (+i \varphi / 2) \\
A_{y} \exp (-i \varphi / 2)
\end{array}\right]
$$

$P_{0} \quad$ Overall Intensity
$P_{0}=I=A_{x}^{2}+A_{y}^{2}$

## Relationship between Jones and Stockes

## Sample Jones Vector

$$
J=\left[\begin{array}{l}
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$$
\begin{aligned}
& P_{0} \quad \text { Overall Intensity } \\
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## Relationship between Jones and Stockes

## Sample Jones Vector

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J=\left[\begin{array}{l}
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A_{y} \exp (-i \varphi / 2)
\end{array}\right]
$$

$P_{0} \quad$ Overall Intensity

$$
\begin{aligned}
& \hline P_{1} \quad \text { Intensity Difference } \\
& \hline
\end{aligned}
$$

$$
\begin{aligned}
& P_{2} \quad \text { in a } \pi / 4 \text { Tilted Basis } \\
& J_{\pi / 4}=\left[\begin{array}{cc}
\cos \left(-\frac{\pi}{4}\right) & -\sin \left(-\frac{\pi}{4}\right) \\
\sin \left(-\frac{\pi}{4}\right) & \cos \left(-\frac{\pi}{4}\right)
\end{array}\right] J
\end{aligned}
$$

## Relationship between Jones and Stockes

## Sample Jones Vector

$$
J=\left[\begin{array}{l}
A_{x} \exp (+i \varphi / 2) \\
A_{y} \exp (-i \varphi / 2)
\end{array}\right]
$$

| $P_{0}$ | Overall Intensity |
| :---: | :---: |
|  | $P_{0}=I=A_{x}^{2}+A_{y}^{2}$ |
| $P_{2}$ | in a $\pi / 4$ Tilted Basis |
|  | $J_{\pi / 4}=\frac{\sqrt{2}}{2}\left[\begin{array}{cc}1 & 1 \\ -1 & 1\end{array}\right] J$ |

$P_{1} \quad$ Intensity Difference

$$
P_{1}=I_{x}-I_{y}=A_{x}^{2}-A_{y}^{2}
$$

Polarization states

## Relationship between Jones and Stockes

## Sample Jones Vector

$$
J=\left[\begin{array}{l}
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## Relationship between Jones and Stockes

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\end{array}\right]
$$

$P_{0} \quad$ Overall Intensity

$P_{0}=I=A_{x}^{2}+A_{y}^{2}$
$P_{2}$ in a $\pi / 4$ Tilted Basis
$J_{\pi / 4}=$
$\sqrt{2}\left[\begin{array}{c}A_{x} e^{+i \varphi / 2}+A_{y} e^{-i \varphi / 2} \\ -A_{x} e^{+i \varphi / 2}+A_{y} e^{-i \varphi / 2}\end{array}\right]$
$P_{2}=\overline{J_{\pi / 4}^{x}} \cdot J_{\pi / 4}^{x}-\overline{J_{\pi / 4}^{y}} \cdot J_{\pi / 4}^{y}=$
$2 A_{x} A_{y} \cos (\varphi)$

## Relationship between Jones and Stockes

## Sample Jones Vector

$$
J=\left[\begin{array}{l}
A_{x} \exp (+i \varphi / 2) \\
A_{y} \exp (-i \varphi / 2)
\end{array}\right]
$$

$P_{0} \quad$ Overall Intensity
$P_{2}$
in a $\pi / 4$ Tilted Basis
$J_{\pi / 4}=$
$\sqrt{2}\left[\begin{array}{c}A_{x} e^{+i \varphi / 2}+A_{y} e^{-i \varphi / 2} \\ -A_{x} e^{+i \varphi / 2}+A_{y} e^{-i \varphi / 2}\end{array}\right]$
$P_{2}=\overline{J_{\pi / 2}^{x}} \cdot J_{\pi / 4}^{x}-\overline{J_{\pi / 4}^{y}} \cdot J_{\pi / 4}^{y}=$
$2 A_{x} A_{y} \cos (\varphi)$

$$
\begin{aligned}
& \text { P1 Intensity Difference } \\
& \\
& P_{1}=I_{x}-I_{y}=A_{x}^{2}-A_{y}^{2}
\end{aligned}
$$

$$
\begin{gathered}
P_{3} \quad \text { in a Circular Basis } \\
J_{\text {Cir }}=\frac{1}{\sqrt{2}}\left[\begin{array}{c}
A_{x} e^{+i \varphi / 2}-i A_{y} e^{-i \varphi / 2} \\
A_{x} e^{+i \varphi / 2}+i A_{y} e^{-i \varphi / 2}
\end{array}\right] \\
P_{3}=\overline{J_{\mathrm{Cir}}^{x}} \cdot J_{\mathrm{Cir}}^{x}-\overline{J_{\text {Cir }}^{y}} \cdot J_{\mathrm{Cir}}^{y}= \\
2 A_{x} A_{y} \sin (\varphi)
\end{gathered}
$$

## Relationship between Jones and Stockes

## Sample Jones Vector

$$
J=\left[\begin{array}{l}
A_{x} \exp (+i \varphi / 2) \\
A_{y} \exp (-i \varphi / 2)
\end{array}\right]
$$

## 4 dependent parameters

$$
P_{0}^{2}=P_{1}^{2}+P_{2}^{2}+P_{3}^{2}
$$

$P_{0} \quad$ Overall Intensity

$$
P_{0}=I=A_{x}^{2}+A_{y}^{2}
$$

$P_{2} \quad$ in a $\pi / 4$ Tilted Basis

$$
\begin{gathered}
J_{\pi / 4}= \\
\sqrt{2}\left[\begin{array}{c}
A_{x} e^{+i \varphi / 2}+A_{y} e^{-i \varphi / 2} \\
-A_{x} e^{+i \varphi / 2}+A_{y} e^{-i \varphi / 2}
\end{array}\right] \\
P_{2}=\overline{J_{\pi / 4}^{x}} \cdot J_{\pi / 4}^{x}-\overline{J_{\pi / 4}^{y}} \cdot J_{\pi / 4}^{y}= \\
2 A_{x} A_{y} \cos (\varphi)
\end{gathered}
$$

## The Poincare Sphere

Polarization states can be described geometrically on a sphere

## Recall the Stokes parameters

$$
P_{0}^{2}=P_{1}^{2}+P_{2}^{2}+P_{3}^{2}
$$

## Normalized Stokes parameters

## The Poincare Sphere

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## Recall the Stokes parameters

$$
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$$

## Normalized Stokes parameters

$$
S_{i}=P_{i} / P_{0}
$$

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$$

## Normalized Stokes parameters

$$
S_{i}=P_{i} / P_{0}
$$

## Unit Radius Sphere

$$
\sum_{i=1}^{3} S_{i}^{2}=1
$$

Polarization states
Jones Calculus
Stokes parameters and the Poincare Sphere

## The Poincare Sphere

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\sum_{i=1}^{3} S_{i}^{2}=1
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## ( $S_{1}, S_{2}, S_{3}$ ) on a unit radius sphere



Figures from [Hua94]

Polarization states
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## The Poincare Sphere

Polarization states can be described geometrically on a sphere

Normalized Stokes parameters

$$
S_{i}=P_{i} / P_{0}
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## Unit Radius Sphere

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\sum_{i=1}^{3} S_{i}^{2}=1
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## General Polarisation



## $\left(S_{1}, S_{2}, S_{3}\right)$ on a unit radius sphere

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Figures from [Hua94]

Polarization Optics, version 2.0, frame 14
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三

## A polarizer lets one component through

Polarizer aligned with $x$ : its action on two orthogonal polarizations

- Lets through the linear polarization along $x:\left[\begin{array}{l}1 \\ 0\end{array}\right] \longrightarrow\left[\begin{array}{l}1 \\ 0\end{array}\right]$ - Blocks the linear polarization along $y:\left[\begin{array}{l}0 \\ 1\end{array}\right] \longrightarrow\left[\begin{array}{l}0 \\ 0\end{array}\right]$


## A polarizer lets one component through

Polarizer aligned with $x$ : its action on two orthogonal polarizations

- Lets through the linear polarization along $x:\left[\begin{array}{l}1 \\ 0\end{array}\right] \longrightarrow\left[\begin{array}{l}1 \\ 0\end{array}\right]$
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## A polarizer lets one component through

Polarizer aligned with $x$ : its action on two orthogonal polarizations

- Lets through the linear polarization along $x:\left[\begin{array}{l}1 \\ 0\end{array}\right] \longrightarrow\left[\begin{array}{l}1 \\ 0\end{array}\right]$
- Blocks the linear polarization along $y:\left[\begin{array}{l}0 \\ 1\end{array}\right] \longrightarrow\left[\begin{array}{l}0 \\ 0\end{array}\right]$
x polarizer Jones matrix
in this basis

$$
\left[\begin{array}{ll}
1 & 0 \\
0 & 0
\end{array}\right]
$$

## A quarter wave plate adds a $\pi / 2$ phase shift

## Birefringent material: $n_{1}$ along $x$ and $n_{2}$ along $y \quad$ thickness $e$

- Linear polarization along $x$ : phase shift is $k e=k_{0} n_{1} e$
- Linear polarization along $y$ : phase shift is $k e=k_{0} n_{2} e$


## A quarter wave plate adds a $\pi / 2$ phase shift

## Birefringent material: $n_{1}$ along $x$ and $n_{2}$ along $y \quad$ thickness $e$ <br> - Linear polarization along $x$ : phase shift is $k e=k_{0} n_{1} e$ <br> - Linear polarization along $y$ : phase shift is $k e=k_{0} n_{2} e$ <br> Jones matrix $\left[\begin{array}{cc}e^{i k_{0} n_{1} e} & 0 \\ 0 & e^{i k_{0} n_{2} e}\end{array}\right]$ <br> in this basis <br> $$
\left[\begin{array}{cc} e^{i k_{0} n_{1} e} & 0 \\ 0 & e^{i k_{0} n_{2} e} \end{array}\right]
$$

## A quarter wave plate adds a $\pi / 2$ phase shift

$$
\begin{aligned}
& \text { Birefringent material: } n_{1} \text { along } x \text { and } n_{2} \text { along } y \quad \text { thickness } e \\
& \text { - Linear polarization along } x \text { : phase shift is } k e=k_{0} n_{1} e \\
& \text { - Linear polarization along } y \text { : phase shift is } k e=k_{0} n_{2} e \\
& \hline \text { Jones matrix } \\
& \hline\left[\begin{array}{cc}
e^{i k_{0} n_{1} e} & 0 \\
0 & e^{i k_{0} n_{2} e}
\end{array}\right]=e^{i k_{0} n_{1} e}\left[\begin{array}{cc}
1 & 0 \\
0 & e^{i k_{0}\left(n_{2}-n_{1}\right) e}
\end{array}\right]
\end{aligned}
$$

## A quarter wave plate adds a $\pi / 2$ phase shift

Birefringent material: $n_{1}$ along $x$ and $n_{2}$ along $y$

- Linear polarization along $x$ : phase shift is $k e=k_{0} n_{1} e$
- Linear polarization along $y$ : phase shift is $k e=k_{0} n_{2} e$
Jones matrix
$\left[\begin{array}{cc}e^{i k_{0} n_{1} e} & 0 \\ 0 & e^{i k_{0} n_{2} e}\end{array}\right]=e^{i k_{0} n_{1} e}\left[\begin{array}{cc}1 & 0 \\ 0 & e^{i k_{0}\left(n_{2}-n_{1}\right) e}\end{array}\right]=e^{i k_{0} n_{1} e}\left[\begin{array}{cc}1 & 0 \\ 0 & e^{ \pm i \frac{\pi}{2}}\end{array}\right]$


## A quarter wave plate adds a $\pi / 2$ phase shift

Birefringent material: $n_{1}$ along $x$ and $n_{2}$ along $y$

- Linear polarization along $x$ : phase shift is $k e=k_{0} n_{1} e$
- Linear polarization along $y$ : phase shift is $k e=k_{0} n_{2} e$

| Jones matrix | in this basis |
| ---: | :--- |
| $\left[\begin{array}{cc}e^{i k_{0} n_{1} e} & 0 \\ 0 & e^{i k_{0} n_{2} e}\end{array}\right]=e^{i k_{0} n_{1} e}\left[\begin{array}{cc}1 & 0 \\ 0 & e^{ \pm i \frac{\pi}{2}}\end{array}\right]=e^{i k_{0} n_{1} e}\left[\begin{array}{cc}1 & 0 \\ 0 & \pm i\end{array}\right]$ |  |

## A quarter wave plate adds a $\pi / 2$ phase shift

> Birefringent material: $n_{1}$ along $x$ and $n_{2}$ along $y$
> - Linear polarization along $x$ : phase shift is $k e=k_{0} n_{1} e$
> - Linear polarization along $y$ : phase shift is $k e=k_{0} n_{2} e$

## A quarter wave plate adds a $\pi / 2$ phase shift

## Birefringent material: $n_{1}$ along $x$ and $n_{2}$ along $y$

- Linear polarization along $x$ : phase shift is $k e=k_{0} n_{1} e$
- Linear polarization along $y$ : phase shift is $k e=k_{0} n_{2} e$
Jones matrix in this basis

$$
\left[\begin{array}{cc}
e^{i k_{0} n_{1} e} & 0 \\
0 & e^{i k_{0} n_{2} e}
\end{array}\right] \approx\left[\begin{array}{cc}
1 & 0 \\
0 & \pm i
\end{array}\right]
$$

## Homework

Think about Half Wave Plates
(1) The physics of polarization optics

- Polarization states
- Jones Calculus
- Stokes parameters and the Poincare Sphere
(2) Polarized light propagation
- Jones Matrices Examples
- Matrix, basis \& eigen polarizations
- Jones Matrices Composition
(3) Partially polarized light
- Formalisms used
- Propagation through optical devices


## Eigen Polarizations

Eigen polarization are polarizations that do not change upon propagation

## Eigen Vectors $\lambda \in \mathbb{C}$ <br> $M \cdot v=\lambda v \Leftrightarrow \begin{aligned} & v \text { is an eigen vector } \\ & \lambda \text { is its eigen value }\end{aligned}$

๑) $Q \curvearrowright$

## Eigen Polarizations

Eigen polarization are polarizations that do not change upon propagation

| Eigen Vectors | $\lambda \in \mathbb{C}$ |
| :--- | :--- |
| $M \cdot v=\lambda v \Leftrightarrow$$v$ is an eigen vector <br> $\lambda$ is its eigen value | $\bullet J$ and $\lambda J$ describe the same |
| polarization |  |

## Polarization unchanged

- $J$ and $\lambda J$ describe the same polarization
- Intensity changes


## Eigen Polarizations

Eigen polarization are polarizations that do not change upon propagation

| Eigen Vectors | $\lambda \in \mathbb{C}$ |
| :--- | :--- |
| $M \cdot v=\lambda v \Leftrightarrow$$v$ is an eigen vector <br> $\lambda$ is its eigen value |  |

## Handy basis

A matrix is diagonal in its eigen basis

- Polarizer eigen basis is along its axes
- Bi-refringent plate eigen basis is along its axes

Find the eigen polarizations for an optically active material that rotates any linear polarisation by an angle $\phi$

## Eigen Polarizations

Eigen polarization are polarizations that do not change upon propagation

| Eigen Vectors | $\lambda \in \mathbb{C}$ |
| :--- | ---: |
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- Polarizer eigen basis is along its axes
- Bi-refringent plate eigen basis is along its axes
๑) $\propto$


## Eigen Polarizations

Eigen polarization are polarizations that do not change upon propagation

$$
\begin{aligned}
& \text { Eigen Vectors } \quad \lambda \in \mathbb{C} \\
& M \cdot v=\lambda v \Leftrightarrow \begin{array}{l}
v \text { is an eigen vector } \\
\lambda \text { is its eigen value }
\end{array}
\end{aligned}
$$

## Polarization unchanged

- $J$ and $\lambda J$ describe the same polarization
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## Handy basis

A matrix is diagonal in its eigen basis

- Polarizer eigen basis is along its axes
- Bi-refringent plate eigen basis is along its axes


## Homework

Find the eigen polarizations for an optically active material that rotates any linear polarisation by an angle $\phi$

## A polarizer in a rotated basis

## In its eigen basis

Eigen basis Jones matrix : $\mathbf{P}_{\mathbf{x}}=\left[\begin{array}{ll}1 & 0 \\ 0 & 0\end{array}\right]$


## A polarizer in a rotated basis

## In its eigen basis

Eigen basis Jones matrix : $\mathbf{P}_{\mathbf{x}}=\left[\begin{array}{ll}1 & 0 \\ 0 & 0\end{array}\right]$
$\square$
Change base through $-\theta$ rotation Transformation Matrix

## A polarizer in a rotated basis

## In its eigen basis

Eigen basis Jones matrix : $\mathbf{P}_{\mathbf{x}}=\left[\begin{array}{ll}1 & 0 \\ 0 & 0\end{array}\right]$
When transmitted polarization is $\theta$ tilted
Change base through - $\theta$ rotation Transformation Matrix

$$
\mathbf{R}(\theta)=\left[\begin{array}{cc}
\cos (\theta) & -\sin (\theta) \\
\sin (\theta) & \cos (\theta)
\end{array}\right]
$$

## A polarizer in a rotated basis

## In its eigen basis

Eigen basis Jones matrix : $\mathbf{P}_{\mathbf{x}}=\left[\begin{array}{ll}1 & 0 \\ 0 & 0\end{array}\right]$
When transmitted polarization is $\theta$ tilted
Change base through - $\theta$ rotation Transformation Matrix

$$
\begin{gathered}
\mathbf{R}(\theta)=\left[\begin{array}{cc}
\cos (\theta) & -\sin (\theta) \\
\sin (\theta) & \cos (\theta)
\end{array}\right] \\
\mathbf{P}(\theta)=\mathbf{R}(-\theta)^{-1}\left[\begin{array}{cc}
1 & 0 \\
0 & 0
\end{array}\right] \mathbf{R}(-\theta)
\end{gathered}
$$

## A polarizer in a rotated basis

## In its eigen basis

Eigen basis Jones matrix : $\mathbf{P}_{\mathbf{x}}=\left[\begin{array}{ll}1 & 0 \\ 0 & 0\end{array}\right]$
When transmitted polarization is $\theta$ tilted
Change base through - $\theta$ rotation Transformation Matrix

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& \mathbf{R}(\theta)=\left[\begin{array}{cc}
\cos (\theta) & -\sin (\theta) \\
\sin (\theta) & \cos (\theta)
\end{array}\right] \\
& \mathbf{P}(\theta)=\mathbf{R}(\theta)\left[\begin{array}{ll}
1 & 0 \\
0 & 0
\end{array}\right] \mathbf{R}(-\theta)
\end{aligned}
$$

## A polarizer in a rotated basis

## In its eigen basis

Eigen basis Jones matrix : $\mathbf{P}_{\mathbf{x}}=\left[\begin{array}{ll}1 & 0 \\ 0 & 0\end{array}\right]$
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\cos (\theta) & -\sin (\theta) \\
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\end{array}\right]
$$

$$
\mathbf{P}(\theta)=\mathbf{R}(\theta)\left[\begin{array}{ll}
1 & 0 \\
0 & 0
\end{array}\right] \mathbf{R}(-\theta)=\left[\begin{array}{cc}
\cos ^{2}(\theta) & \sin (\theta) \cos (\theta) \\
\sin (\theta) \cos (\theta) & \sin ^{2}(\theta)
\end{array}\right]
$$

## A polarizer in a rotated basis

## When transmitted polarization is $\theta$ tilted

Change base through - $\theta$ rotation Transformation Matrix

$$
\mathbf{R}(\theta)=\left[\begin{array}{cc}
\cos (\theta) & -\sin (\theta) \\
\sin (\theta) & \cos (\theta)
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$$

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1 & 0 \\
0 & 0
\end{array}\right] \mathbf{R}(-\theta)=\left[\begin{array}{cc}
\cos ^{2}(\theta) & \sin (\theta) \cos (\theta) \\
\sin (\theta) \cos (\theta) & \sin ^{2}(\theta)
\end{array}\right]
$$

## Homework

- Find the same result with physics arguments only
- Find the power transmitted through such a polarizer from a $x$ polarized plane wave


## Changing basis in the general case

## Using the Transformation Matrix

- If basis $B_{1}$ is deduded from basis $B_{0}$ by transformation $P$ : $B_{1}=P B_{0}$
- Jones Matrix is transformed using $J_{1}=P^{-1} J_{0} P$


## Changing basis in the general case

## Using the Transformation Matrix

- If basis $B_{1}$ is deduded from basis $B_{0}$ by transformation $P$ :

$$
B_{1}=P B_{0}
$$

- Jones Matrix is transformed using $J_{1}=P^{-1} J_{0} P$

From linear to circular

- Optically Active media in a linear basis :

$$
J=\left[\begin{array}{cc}
\cos (\phi) & \sin (\phi) \\
-\sin (\phi) & \cos (\phi)
\end{array}\right]
$$

- Transformation Matrix to a circular basis $P=\left[\begin{array}{cc}1 & 1 \\ 1 & -i\end{array}\right]$
- $P^{-1} M P=\left[\begin{array}{cc}e^{i \phi} & 0 \\ 0 & e^{-i \phi}\end{array}\right]$


## Anisotropy can be linear and circular

## Linear Anisotropy

- Orthogonal eigen linear polarizations
- Different index $n_{1}$ \& $n_{2}$

Eigen Jones Matrix

$$
\left[\begin{array}{cc}
1 & 0 \\
0 & e^{i \theta}
\end{array}\right]
$$

Orthogonal linear polarisations basis

## Anisotropy can be linear and circular

## Linear Anisotropy

- Orthogonal eigen linear polarizations
- Different index $n_{1} \& n_{2}$

Eigen Jones Matrix

$$
\left[\begin{array}{cc}
1 & 0 \\
0 & e^{i \theta}
\end{array}\right]
$$

Orthogonal linear polarisations basis

## Circular Anisotropy

- Orthogonal eigen Circular polarizations
- Different index $n_{1}$ \& $n_{2}$

Eigen Jones Matrix

$$
\left[\begin{array}{cc}
1 & 0 \\
0 & e^{i \theta}
\end{array}\right]
$$

Orthogonal Circular basis

## Anisotropy can be linear and circular

## Linear Anisotropy

- Orthogonal eigen linear polarizations
- Different index $n_{1}$ \& $n_{2}$

Eigen Jones Matrix

$$
\left[\begin{array}{cc}
1 & 0 \\
0 & e^{i \theta}
\end{array}\right]
$$

Orthogonal linear polarisations basis

## Circular Anisotropy

- Orthogonal eigen Circular polarizations
- Different index $n_{1}$ \& $n_{2}$

Eigen Jones Matrix
$\left[\begin{array}{cc}1 & 0 \\ 0 & e^{i \theta}\end{array}\right]$

Orthogonal Circular basis
Back to linear basis

$$
\begin{array}{ll}
{\left[\begin{array}{cc}
\cos \left(\frac{\theta}{2}\right) & \sin \left(\frac{\theta}{2}\right) \\
-\sin \left(\frac{\theta}{2}\right) & \cos \left(\frac{\theta}{2}\right)
\end{array}\right]} \\
\text { Optically Active media }
\end{array}
$$

(1) The physics of polarization optics

- Polarization states
- Jones Calculus
- Stokes parameters and the Poincare Sphere
(2) Polarized light propagation
- Jones Matrices Examples
- Matrix, basis \& eigen polarizations
- Jones Matrices Composition
(3) Partially polarized light
- Formalisms used
- Propagation through optical devices


## Jones Matrices Composition

The Jones matrices of cascaded optical elements can be composed through Matrix multiplication

## Matrix composition

- If a $\overrightarrow{J_{0}}$ incident light passes through $\mathbf{M}_{\mathbf{1}}$ and $\mathbf{M}_{\mathbf{2}}$ in that order
- First transmission: $\mathbf{M}_{\mathbf{1}} \overrightarrow{J_{0}}$
- Second transmission: $\mathbf{M}_{2} \mathbf{M}_{\mathbf{1}} \overrightarrow{J_{0}}$
- Composed Jones Matrix: $\mathbf{M}_{\mathbf{2}} \mathbf{M}_{\mathbf{1}}$


## Jones Matrices Composition

The Jones matrices of cascaded optical elements can be composed through Matrix multiplication

## Matrix composition

- If a $\overrightarrow{J_{0}}$ incident light passes through $\mathbf{M}_{\mathbf{1}}$ and $\mathbf{M}_{\mathbf{2}}$ in that order
- First transmission: $\mathbf{M}_{\mathbf{1}} \overrightarrow{\mathrm{J}_{0}}$
- Second transmission: $\mathbf{M}_{2} \mathbf{M}_{\mathbf{1}} \overrightarrow{J_{0}}$
- Composed Jones Matrix : $\mathbf{M}_{\mathbf{2}} \mathbf{M}_{\mathbf{1}}$

Reversed order

## Beware of non commutativity

- Matrix product does not commute in general
- Think of the case of a linear anisotropy followed by optical activity
- in that order
- in the reverse order
(1) The physics of polarization optics
- Polarization states
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三

## Stokes parameters for partially polarized light

 Generalize the coherent definition using the statistical average intensity
## Stokes Vector

$$
\vec{S}=\left[\begin{array}{c}
P_{0} \\
P_{1} \\
P_{2} \\
P_{3}
\end{array}\right]=\left[\begin{array}{c}
\left\langle I_{x}+I_{y}\right\rangle \\
\left\langle I_{x}-I_{y}\right\rangle \\
\left\langle I_{\pi / 4}-I_{-\pi / 4}\right\rangle \\
\left\langle I_{L}-I_{R}\right\rangle
\end{array}\right]
$$



## Stokes parameters for partially polarized light

 Generalize the coherent definition using the statistical average intensity
## Stokes Vector

$$
\vec{S}=\left[\begin{array}{c}
P_{0} \\
P_{1} \\
P_{2} \\
P_{3}
\end{array}\right]=\left[\begin{array}{c}
\left\langle I_{x}+I_{y}\right\rangle \\
\left\langle I_{x}-I_{y}\right\rangle \\
\left\langle I_{\pi / 4}-I_{-\pi / 4}\right\rangle \\
\left\langle I_{L}-I_{R}\right\rangle
\end{array}\right]
$$

## Polarization degree $0 \leq p \leq 1$

$$
p=\frac{\sqrt{P_{1}^{2}+P_{2}^{2}+P_{3}^{2}}}{P_{0}}
$$

Polarized and depolarized sum

## Stokes parameters for partially polarized light

 Generalize the coherent definition using the statistical average intensity
## Stokes Vector

## Polarization degree $0 \leq p \leq 1$

$$
\vec{S}=\left[\begin{array}{c}
P_{0} \\
P_{1} \\
P_{2} \\
P_{3}
\end{array}\right]=\left[\begin{array}{c}
\left\langle I_{x}+I_{y}\right\rangle \\
\left\langle I_{x}-I_{y}\right\rangle \\
\left\langle I_{\pi / 4}-I_{-\pi / 4}\right\rangle \\
\left\langle I_{L}-I_{R}\right\rangle
\end{array}\right]
$$

$$
p=\frac{\sqrt{P_{1}^{2}+P_{2}^{2}+P_{3}^{2}}}{P_{0}}
$$

## Stokes decomposition

$$
\vec{S}=\left[\begin{array}{l}
P_{0} \\
P_{1} \\
P_{2} \\
P_{3}
\end{array}\right]=\left[\begin{array}{c}
p P_{0} \\
P_{1} \\
P_{2} \\
P_{3}
\end{array}\right]+\left[\begin{array}{c}
(1-p) P_{0} \\
0 \\
0 \\
0
\end{array}\right]=\overrightarrow{S_{P}}+\overrightarrow{S_{N P}}
$$

## The Jones Coherence Matrix

## Jones Coherence Matrix

$$
\text { - If } \vec{J}=\left[\begin{array}{l}
A_{x}(t) e^{i \varphi_{x}(t)} \\
A_{y}(t) e^{i \varphi_{y}(t)}
\end{array}\right]
$$

- Meaningless when not monochromatic

Formalisms used

## The Jones Coherence Matrix

## Jones Coherence Matrix

- If $\vec{J}=\left[\begin{array}{l}A_{x}(t) e^{i \varphi_{x}(t)} \\ A_{y}(t) e^{i \varphi_{y}(t)}\end{array}\right]$
- $\Gamma_{i j}=\left\langle\vec{J}_{i}(t) \vec{J}_{j}(t)\right\rangle$


## The Jones Coherence Matrix

## Jones Coherence Matrix

- If $\vec{J}=\left[\begin{array}{l}A_{x}(t) e^{i \varphi_{x}(t)} \\ A_{y}(t) e^{i \varphi_{y}(t)}\end{array}\right]$
- $\Gamma_{i j}=\left\langle\vec{J}_{i}(t) \vec{J}_{j}(t)\right\rangle$
- $\Gamma=\left\langle\overrightarrow{J(t)} \overline{\vec{J}(t)}^{t}\right\rangle$


## The Jones Coherence Matrix

## Jones Vectors are out

- They describe phase differences
- Meaningless when not monochromatic


## Jones Coherence Matrix

- If $\vec{J}=\left[\begin{array}{l}A_{x}(t) e^{i \varphi_{x}(t)} \\ A_{y}(t) e^{i \varphi_{y}(t)}\end{array}\right]$
- $\Gamma_{i j}=\left\langle\vec{J}_{i}(t) \vec{J}_{j}(t)\right\rangle$
- $\Gamma=\left\langle\overrightarrow{J(t)} \overline{\overline{J(t)}^{t}}\right\rangle$

Coherence Matrix: explicit formulation

$$
\Gamma=\left[\begin{array}{ll}
\Gamma_{x x} & \Gamma_{x y} \\
\Gamma_{y x} & \Gamma_{y y}
\end{array}\right]
$$

## The Jones Coherence Matrix

## Jones Coherence Matrix

## Jones Vectors are out

- They describe phase differences
- Meaningless when not monochromatic

$$
\text { - If } \vec{J}=\left[\begin{array}{l}
A_{x}(t) e^{i \varphi_{x}(t)} \\
A_{y}(t) e^{i \varphi_{y}(t)}
\end{array}\right]
$$

- $\Gamma_{i j}=\left\langle\vec{J}_{i}(t) \vec{J}_{j}(t)\right\rangle$
- $\Gamma=\left\langle\overrightarrow{J(t)}{\overrightarrow{J J} \vec{J}^{t}}\right\rangle$


## Coherence Matrix: explicit formulation

$$
\Gamma=\left[\begin{array}{cc}
\left.\left.\langle | A_{x}(t)\right|^{2}\right\rangle & \left\langle A_{x}(t) \overline{A_{y}(t)} e^{i\left(\varphi_{x}-\varphi_{y}\right)}\right\rangle \\
\left\langle\overline{A_{x}(t)} A_{y}(t) e^{-i\left(\varphi_{x}-\varphi_{y}\right)}\right\rangle & \left.\left.\langle | A_{y}(t)\right|^{2}\right\rangle
\end{array}\right]
$$

## Jones Coherence Matrix: properties

## The Coherence Matrix

$$
\Gamma=\left[\begin{array}{cc}
\left.\left.\langle | A_{x}(t)\right|^{2}\right\rangle & \left\langle A_{x}(t) \overline{A_{y}(t)} e^{i\left(\varphi_{x}-\varphi_{y}\right)}\right\rangle \\
\left\langle\overline{A_{x}(t)} A_{y}(t) e^{-i\left(\varphi_{x}-\varphi_{y}\right)}\right\rangle & \left.\left.\langle | A_{y}(t)\right|^{2}\right\rangle
\end{array}\right]
$$

## Trace is Intensity

$$
\operatorname{Tr}(\Gamma)=1
$$

## Jones Coherence Matrix: properties

## The Coherence Matrix

$$
\Gamma=\left[\begin{array}{cc}
\left.\left.\langle | A_{x}(t)\right|^{2}\right\rangle & \left\langle A_{x}(t) \overline{A_{y}(t)} e^{i\left(\varphi_{x}-\varphi_{y}\right)}\right\rangle \\
\left\langle\overline{A_{x}(t)} A_{y}(t) e^{-i\left(\varphi_{x}-\varphi_{y}\right)}\right\rangle & \left.\left.\langle | A_{y}(t)\right|^{2}\right\rangle
\end{array}\right]
$$

## Trace is Intensity

$$
\operatorname{Tr}(\Gamma)=1
$$

Base change
Transformation $\mathbf{P}$

$$
\mathbf{P}^{-1} \Gamma \mathbf{P}
$$

## Jones Coherence Matrix: properties

## The Coherence Matrix

$$
\Gamma=\left[\begin{array}{cc}
\left.\left.\langle | A_{x}(t)\right|^{2}\right\rangle & \left\langle A_{x}(t) \overline{A_{y}(t)} e^{i\left(\varphi_{x}-\varphi_{y}\right)}\right\rangle \\
\left\langle\overline{A_{x}(t)} A_{y}(t) e^{-i\left(\varphi_{x}-\varphi_{y}\right)}\right\rangle & \left.\left.\langle | A_{y}(t)\right|^{2}\right\rangle
\end{array}\right]
$$

Trace is Intensity
Base change Transformation $\mathbf{P}$

$$
\operatorname{Tr}(\Gamma)=1
$$

$$
\mathbf{P}^{-1} \Gamma \mathbf{P}
$$

Relationship with Stokes parameters
from definition

$$
\left[\begin{array}{l}
P_{0} \\
P_{1} \\
P_{2} \\
P_{3}
\end{array}\right]=\left[\begin{array}{cccc}
1 & 1 & 0 & 0 \\
1 & -1 & 0 & 0 \\
0 & 0 & 1 & 1 \\
0 & 0 & -i & i
\end{array}\right]\left[\begin{array}{l}
\Gamma_{x x} \\
\Gamma_{y y} \\
\Gamma_{x y} \\
\Gamma_{y x}
\end{array}\right]
$$

## Jones Coherence Matrix: properties

Trace is Intensity

$$
\operatorname{Tr}(\Gamma)=1
$$

## Base change

Transformation $\mathbf{P}$

## $\mathbf{P}^{-1}$ 「 $\mathbf{P}$

Relationship with Stokes parameters
from definition

$$
\left[\begin{array}{l}
P_{0} \\
P_{1} \\
P_{2} \\
P_{3}
\end{array}\right]=\left[\begin{array}{cccc}
1 & 1 & 0 & 0 \\
1 & -1 & 0 & 0 \\
0 & 0 & 1 & 1 \\
0 & 0 & -i & i
\end{array}\right]\left[\begin{array}{l}
\Gamma_{x x} \\
\Gamma_{y y} \\
\Gamma_{x y} \\
\Gamma_{y x}
\end{array}\right]
$$

Inverse relationship

$$
\left[\begin{array}{l}
\Gamma_{x x} \\
\Gamma_{y y} \\
\Gamma_{x y} \\
\Gamma_{y x}
\end{array}\right]=\frac{1}{2}\left[\begin{array}{cccc}
1 & 1 & 0 & 0 \\
1 & -1 & 0 & 0 \\
0 & 0 & 1 & i \\
0 & 0 & 1 & -i
\end{array}\right]\left[\begin{array}{l}
P_{0} \\
P_{1} \\
P_{2} \\
P_{3}
\end{array}\right]
$$

## Coherence Matrix: further properties

## Polarization degree

$$
p=\sqrt{\frac{P_{1}^{2}+P_{2}^{2}+P_{3}^{2}}{P_{0}^{2}}}=\sqrt{1-\frac{4\left(\Gamma_{x x} \Gamma_{y y}-\Gamma_{x y} \Gamma_{y x}\right)}{\left(\Gamma_{x x}+\Gamma_{y y}\right)^{2}}}=\sqrt{1-\frac{4 \operatorname{Det}(\Gamma)}{\operatorname{Tr}(\Gamma)^{2}}}
$$


parameters

## Coherence Matrix: further properties

## Polarization degree

$$
p=\sqrt{\frac{P_{1}^{2}+P_{2}^{2}+P_{3}^{2}}{P_{0}^{2}}}=\sqrt{1-\frac{4\left(\Gamma_{x x} \Gamma_{y y}-\Gamma_{x y} \Gamma_{y x}\right)}{\left(\Gamma_{x x}+\Gamma_{y y}\right)^{2}}}=\sqrt{1-\frac{4 \operatorname{Det}(\Gamma)}{\operatorname{Tr}(\Gamma)^{2}}}
$$

$\Gamma$ Decomposition in polarized and depolarized components

- $\Gamma=\Gamma_{P}+\Gamma_{N P}$
- Find $\Gamma_{P}$ and $\Gamma_{N P}$ using the relationship with the Stokes parameters
(1) The physics of polarization optics
- Polarization states
- Jones Calculus
- Stokes parameters and the Poincare Sphere
(2) Polarized light propagation
- Jones Matrices Examples
- Matrix, basis \& eigen polarizations
- Jones Matrices Composition
(3) Partially polarized light
- Formalisms used
- Propagation through optical devices


## Propagation of the Coherence Matrix

## Jones Calculus

- If incoming polarization is $\overrightarrow{J(t)}$
- Output one is $\overrightarrow{J^{\prime}(t)}=\mathbf{M} \overrightarrow{J(t)}$


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Coherence Matrix if $\mathbf{M}$ is unitary

- M unitary means : linear and/or circular anisotropy only


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- $\Gamma^{\prime}=\mathbf{M}\left\langle\overrightarrow{J(t)} \overrightarrow{\left.\bar{J}^{\prime} t\right)^{t}}\right\rangle \mathbf{M}^{-1}$

Basis change

## Propagation of the Coherence Matrix

## Jones Calculus

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Basis change
Polarization degree

- Unaltered for unitary operators

Tr and Det are unaltered

## Propagation of the Coherence Matrix

## Jones Calculus

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Basis change
Polarization degree

- Unaltered for unitary operators $\quad$ Tr and Det are unaltered
- Not the case if a polarizer is present: $p$ becomes 1


## Mueller Calculus

Propagating the Jones coherence matrix is difficult if the operator is not unitary

Jones Calculus raises some difficulties

- Coherence matrix OK for partially polarized light
- Propagation through unitary optical devices
(linear or circular anisotropy only)
$\square$
- Hard Times if Polarizers are present


三

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The Stokes parameters may be an alternative

- Describing intensity, they can be readily measurered


## Mueller Calculus

Propagating the Jones coherence matrix is difficult if the operator is not unitary

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- Coherence matrix OK for partially polarized light
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- Hard Times if Polarizers are present

The Stokes parameters may be an alternative

- Describing intensity, they can be readily measurered
- We will show they can be propagated using $4 \times 4$ real matrices
- They are the Mueller matrices


## The projection on a polarization state

Matrix of the polarizer with axis parallel to $\vec{V}$

## Projection on $\vec{V}$ in Jones Basis

- Orthogonal Linear Polarizations Basis: $\vec{X}$ and $\vec{Y}$
- Normed Projection Base Vector :
- $\vec{V}=A_{x} e^{-i \frac{\varphi}{2}} \vec{X}+A_{y} e^{i \frac{\varphi}{2}} \vec{Y}$
- $\overline{\vec{V}}^{\mathrm{t}} \vec{V}=1$


## The projection on a polarization state

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- $\overrightarrow{\vec{V}}^{\mathrm{t}} \vec{V}=1$
- $\mathbf{P}_{\mathbf{V}}=\vec{V} \overline{\vec{V}}^{\mathrm{t}}{ }^{a}$
${ }^{a}$ Easy to check in the projection eigen basis


## The Pauli Matrices

## A base for the 4D $2 \times 2$ matrix vector space

$$
\sigma_{0}=\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right], \sigma_{1}=\left[\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right], \sigma_{2}=\left[\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right], \sigma_{3}=\left[\begin{array}{cc}
0 & -i \\
i & 0
\end{array}\right]
$$

Pv decomposition

$$
\mathbf{P}_{\mathbf{v}}=\frac{1}{2}\left(p_{0} \sigma_{0}+p_{1} \sigma_{1}+p_{2} \sigma_{2}+p_{3} \sigma_{3}\right)
$$

## Pv composition and Trace property

Trace is the eigen values sum

## Projection property

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Projection property

Projection Trace in its eigen basis

- $\mathbf{P}_{\mathbf{V}}$ eigenvalues : 0 \& 1
$\operatorname{Tr}\left(P_{V}\right)=1$


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Trace is the eigen values sum

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Projection Trace in its eigen basis

- $\mathbf{P}_{\mathbf{V}}$ eigenvalues : 0 \& 1
- $\mathbf{P}_{\mathbf{V}} \sigma_{\mathbf{j}}$ eigenvalues : 0 \& $\alpha$
$\alpha \leq 1$
$\operatorname{Tr}\left(P_{V}\right)=1$
$\operatorname{Tr}\left(P_{V} \sigma_{j}\right)=\alpha$


## $\mathbf{P}_{\mathbf{V}}$ composition and Trace property

Trace is the eigen values sum

## Projection property

Projection Trace in its eigen basis

- $\mathbf{P}_{\mathbf{V}}$ eigenvalues : 0 \& 1
- $\mathbf{P}_{\mathbf{V}} \sigma_{\mathbf{j}}$ eigenvalues : 0 \& $\alpha \quad \alpha \leq 1$
- $\mathbf{P}_{\mathbf{v}} \sigma_{\mathbf{j}}$ eigenvectors are the same as $\mathbf{P}_{\mathbf{V}}$ :
- $\vec{V}$ associated to eigenvalue $\alpha$


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- Project the projection


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- Project the projection


## $\mathbf{P}_{\mathbf{V}}$ Pauli components and physical meaning

Express $p_{i}$ as a function of $\vec{V}$ and the Pauli matrices, then find their signification

$$
\begin{array}{cc}
\overline{\vec{V}^{\mathrm{t}}} \cdot \sigma_{\mathrm{j}} \vec{V}=\operatorname{Tr}\left(P_{V} \sigma_{j}\right) & \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right)=2 \delta_{i j} \\
\operatorname{Tr}\left(P_{V} \sigma_{j}\right)=\frac{1}{2} \sum_{i} \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right) p_{i} &
\end{array}
$$

## Pv Pauli components and physical meaning

Express $p_{i}$ as a function of $\vec{V}$ and the Pauli matrices, then find their signification

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\begin{gathered}
\overline{\overline{\mathrm{V}}^{\mathrm{t}}} \cdot \sigma_{\mathrm{j}} \vec{V}=\operatorname{Tr}\left(P_{V} \sigma_{j}\right) \quad \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right)=2 \delta_{i j} \\
\operatorname{Tr}\left(P_{V} \sigma_{j}\right)=\frac{1}{2} \sum_{i} \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right) p_{i}=\frac{1}{2} \sum_{i} 2 \delta_{i j} p_{i}
\end{gathered}
$$

## $\mathbf{P}_{\mathbf{V}}$ Pauli components and physical meaning

Express $p_{i}$ as a function of $\vec{V}$ and the Pauli matrices, then find their signification

$$
\begin{aligned}
& \vec{V}^{t} \cdot \sigma_{\mathrm{j}} \vec{V}=\operatorname{Tr}\left(P_{V} \sigma_{j}\right) \quad \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right)=2 \delta_{i j} \\
& \overline{\vec{V}^{t}} \cdot \sigma_{\mathbf{j}} \vec{V}=\operatorname{Tr}\left(P_{V} \sigma_{j}\right)=\frac{1}{2} \sum_{i} \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right) p_{i}=\frac{1}{2} \sum_{i} 2 \delta_{i j} p_{i}=p_{j}
\end{aligned}
$$

## $\mathbf{P}_{\mathbf{V}}$ Pauli components and physical meaning

Express $p_{i}$ as a function of $\vec{V}$ and the Pauli matrices, then find their signification

$$
\begin{aligned}
& {\overrightarrow{\vec{V}^{\mathrm{t}}} \cdot \sigma_{\mathrm{j}} \vec{V}=\operatorname{Tr}\left(P_{V} \sigma_{j}\right) \quad \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right)=2 \delta_{i j}}_{\overrightarrow{\vec{V}}^{\mathrm{t}} \cdot \sigma_{\mathrm{j}} \vec{V}=\operatorname{Tr}\left(P_{V} \sigma_{j}\right)=\frac{1}{2} \sum_{i} \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right) p_{i}=\frac{1}{2} \sum_{i} 2 \delta_{i j} p_{i}=p_{j}} .
\end{aligned}
$$

Project the base vectors on $\vec{V}$

- Using $\vec{V}=A_{x} e^{-i \frac{\varphi}{2}} \vec{X}+A_{y} e^{i \frac{\varphi}{2}} \vec{Y}$
- $\operatorname{Prv}_{\mathrm{v}} \vec{X}=A_{x}^{2} \vec{X}+A_{x} A_{y} e^{i \varphi} \vec{Y}$
- $\mathbf{P}_{v} \vec{Y}=A_{y}^{2} \vec{Y}+A_{x} A_{y} e^{-i \varphi} \vec{X}$


## - Using the $\mathrm{P}_{\mathrm{V}}$ decomposition on the Pauli Basis

## Pv Pauli components and physical meaning

Express $p_{i}$ as a function of $\vec{V}$ and the Pauli matrices, then find their signification

$$
\begin{aligned}
& {\overrightarrow{\vec{V}^{\mathrm{t}}} \cdot \sigma_{\mathrm{j}} \vec{V}=\operatorname{Tr}\left(P_{V} \sigma_{j}\right) \quad \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right)=2 \delta_{i j}}_{\overrightarrow{\vec{V}}^{\mathrm{t}} \cdot \sigma_{\mathrm{j}} \vec{V}=\operatorname{Tr}\left(P_{V} \sigma_{j}\right)=\frac{1}{2} \sum_{i} \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right) p_{i}=\frac{1}{2} \sum_{i} 2 \delta_{i j} p_{i}=p_{j}}
\end{aligned}
$$

Project the base vectors on $\vec{V}$

- Using $\vec{V}=A_{x} e^{-i \frac{\varphi}{2}} \vec{X}+A_{y} e^{i \frac{\varphi}{2}} \vec{Y}$
- $\operatorname{Prv}_{\mathrm{v}} \vec{X}=A_{x}^{2} \vec{X}+A_{x} A_{y} e^{i \varphi} \vec{Y}$
- $\mathrm{P}_{v} \vec{Y}=A_{y}^{2} \vec{Y}+A_{x} A_{y} e^{-i \varphi} \vec{X}$
- Using the $\mathbf{P}_{\mathbf{V}}$ decomposition on the Pauli Basis
- $\boldsymbol{P}_{\mathbf{v}} \vec{X}=\frac{1}{2}\left(p_{0}+p_{1}\right) \vec{X}+\frac{1}{2}\left(p_{2}+i p_{3}\right) \vec{Y}$
- $\mathbf{P}_{\mathbf{v}} \vec{Y}=\frac{1}{2}\left(p_{0}-p_{1}\right) \vec{Y}+\frac{1}{2}\left(p_{2}-i p_{3}\right) \vec{X}$


## $\mathbf{P}_{\mathbf{V}}$ Pauli components and physical meaning

Express $p_{i}$ as a function of $\vec{V}$ and the Pauli matrices, then find their signification

$$
\begin{aligned}
& {\overrightarrow{\vec{V}^{\mathrm{t}}} \cdot \sigma_{\mathrm{j}} \vec{V}=\operatorname{Tr}\left(P_{V} \sigma_{j}\right) \quad \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right)=2 \delta_{i j}}_{\overrightarrow{\vec{V}}^{\mathrm{t}} \cdot \sigma_{\mathrm{j}} \vec{V}=\operatorname{Tr}\left(P_{V} \sigma_{j}\right)=\frac{1}{2} \sum_{i} \operatorname{Tr}\left(\sigma_{i} \sigma_{j}\right) p_{i}=\frac{1}{2} \sum_{i} 2 \delta_{i j} p_{i}=p_{j}}
\end{aligned}
$$

Project the base vectors on $\vec{V}$

- Using $\vec{V}=A_{x} e^{-i \frac{\varphi}{2}} \vec{X}+A_{y} e^{i \frac{\varphi}{2}} \vec{Y}$
- $\operatorname{Prv}_{\mathrm{v}} \vec{X}=A_{x}^{2} \vec{X}+A_{x} A_{y} e^{i \varphi} \vec{Y}$
- $\mathrm{P}_{v} \vec{Y}=A_{y}^{2} \vec{Y}+A_{x} A_{y} e^{-i \varphi} \vec{X}$
- Using the $\mathbf{P}_{\mathbf{V}}$ decomposition on the Pauli Basis
- $\boldsymbol{P}_{\mathbf{v}} \vec{X}=\frac{1}{2}\left(p_{0}+p_{1}\right) \vec{X}+\frac{1}{2}\left(p_{2}+i p_{3}\right) \vec{Y}$
- $\mathbf{P}_{\mathbf{v}} \vec{Y}=\frac{1}{2}\left(p_{0}-p_{1}\right) \vec{Y}+\frac{1}{2}\left(p_{2}-i p_{3}\right) \vec{X}$
- Identify


## Pv Pauli composition and Stokes parameters

## Stokes parameters as $\mathrm{Pv}_{\mathrm{V}}$ decomposition on the Pauli base

- $p_{0}=P_{0}=A_{x}^{2}-A_{y}^{2}=I_{x}-I_{y}$
- $p_{1}=P_{1}=A_{x}^{2}-A_{y}^{2}=I_{x}-I_{y}$
- $p_{2}=P_{2}=2 A_{x} A_{y} \cos (\varphi)=I_{\pi / 4}-I_{-\pi / 4}$
- $p_{3}=P_{3}=2 A_{x} A_{y} \sin (\varphi)=I_{\mathrm{L}}-I_{\mathrm{R}}$


## Propagating through devices: Mueller matrices $\overrightarrow{v^{\prime}}=M_{\lrcorner} \vec{V}$

Projection on $\overrightarrow{V^{\prime}}$

$$
\mathbf{P}_{\mathbf{V}^{\prime}}=\overrightarrow{V^{\prime}} \overrightarrow{\bar{V}^{\prime}}
$$

## Propagating through devices: Mueller matrices $\overrightarrow{v^{\prime}}=M_{\lrcorner} \vec{V}$

Projection on $\overrightarrow{V^{\prime}}$

$$
\mathbf{P}_{\mathbf{V}^{\prime}}=\overrightarrow{V^{\prime}} \overline{\vec{V}^{\prime}}=\mathbf{M}_{\mathbf{\jmath}} \vec{V} \overrightarrow{\vec{V}^{t}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}}
$$

## Propagating through devices: Mueller matrices $\overrightarrow{v^{\prime}}=M_{\lrcorner} \vec{V}$

Projection on $\overrightarrow{V^{\prime}}$

$$
\mathbf{P}_{\mathbf{V}^{\prime}}=\overrightarrow{V^{\prime}} \overrightarrow{{\overrightarrow{V^{\prime}}}^{t}}=\mathbf{M}_{\mathbf{J}} \vec{V} \overline{\vec{V}^{t}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}}=\mathbf{M}_{\mathbf{J}} \mathbf{P}_{\mathbf{V}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}}
$$

## Propagating through devices: Mueller matrices $\overrightarrow{V^{\prime}}=M_{\lrcorner} \vec{V}$

Projection on $\overrightarrow{V^{\prime}}$

$$
\mathbf{P}_{\mathbf{V}^{\prime}}=\overrightarrow{V^{\prime}} \overrightarrow{\vec{V}^{\prime}}=\mathbf{M}_{\mathbf{J}} \overrightarrow{V_{V}} \overline{\vec{V}^{\mathrm{t}}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}}=\mathbf{M}_{\mathbf{J}} \mathbf{P}_{\mathbf{V}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}}
$$

Trace relationship

$$
P_{i}^{\prime}=\operatorname{Tr}\left(\mathbf{P}_{\mathbf{v}^{\prime}} \sigma_{i}\right)
$$

## Propagating through devices: Mueller matrices $\overrightarrow{v^{\prime}}=M_{J} \vec{V}$

Projection on $\overrightarrow{V^{\prime}}$

$$
\mathbf{P}_{\mathbf{V}^{\prime}}=\overrightarrow{V^{\prime}} \overline{\bar{V}^{\mathrm{t}}}=\mathbf{M}_{\mathbf{J}} \vec{V} \overline{\vec{V}^{\mathrm{t}}} \overline{\mathbf{M}_{\mathbf{J}}{ }^{\mathrm{t}}}=\mathbf{M}_{\mathbf{J}} \mathbf{P}_{\mathbf{V}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}}
$$

## Trace relationship

$$
P_{i}^{\prime}=\operatorname{Tr}\left(\mathbf{P}_{\mathbf{v}^{\prime}} \sigma_{i}\right)=\operatorname{Tr}\left(\mathbf{M}_{\mathbf{J}} \mathbf{P}_{\mathbf{V}} \overline{\left.\mathbf{M}_{\mathbf{J}}{ }^{\mathrm{t}} \sigma_{\mathbf{i}}\right)}\right.
$$

## Propagating through devices: Mueller matrices $\overrightarrow{v^{\prime}}=M_{J} \vec{V}$

Projection on $\overrightarrow{V^{\prime}}$

$$
\mathbf{P}_{\mathbf{V}^{\prime}}=\overrightarrow{V^{\prime}} \overline{{\overline{V^{\prime}}}^{t}}=\mathbf{M}_{\mathbf{J}} \overrightarrow{\vec{V}^{\prime}} \overline{\vec{V}^{t}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}}=\mathbf{M}_{\mathbf{J}} \mathbf{P}_{\mathbf{V}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}}
$$

Trace relationship

$$
\begin{gathered}
P_{i}^{\prime}=\operatorname{Tr}\left(\mathbf{P}_{\mathbf{V}^{\prime}} \sigma_{i}\right)=\operatorname{Tr}\left(\mathbf{M}_{\mathbf{J}} \mathbf{P}_{\mathbf{V}} \overline{\left.\mathbf{M}_{\mathbf{J}}{ }^{\mathrm{t}} \sigma_{\mathbf{i}}\right)=}\right. \\
\frac{1}{2} \sum_{j=0}^{3} \operatorname{Tr}\left(\mathbf{M}_{\mathbf{J}} \sigma_{\mathbf{j}} \overline{\left.\mathbf{M}_{\mathbf{J}}{ }^{\mathrm{t}} \sigma_{\mathbf{i}}\right) P_{j}}\right.
\end{gathered}
$$

## Propagating through devices: Mueller matrices $\overrightarrow{v^{\prime}}=M_{\jmath} \vec{V}$

Projection on $\overrightarrow{V^{\prime}}$

$$
\mathbf{P}_{\mathbf{V}^{\prime}}=\overrightarrow{V^{\prime}} \overline{{\overline{V^{\prime}}}^{t}}=\mathbf{M}_{\mathbf{J}} \overrightarrow{\vec{V}^{\prime}} \overline{\vec{V}^{t}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}}=\mathbf{M}_{\mathbf{J}} \mathbf{P}_{\mathbf{V}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}}
$$

Trace relationship

$$
\begin{gathered}
P_{i}^{\prime}=\operatorname{Tr}\left(\mathbf{P}_{\mathbf{V}^{\prime}} \sigma_{i}\right)=\operatorname{Tr}\left(\mathbf{M}_{\mathbf{J}} \mathbf{P}_{\mathbf{V}} \overline{\left.\mathbf{M}_{\mathbf{J}}{ }^{\mathrm{t}} \sigma_{\mathbf{i}}\right)=}\right. \\
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\end{gathered}
$$

Mueller matrix

$$
\overrightarrow{S^{\prime}}=\mathbf{M}_{\mathbf{M}} \vec{S}
$$

$$
\left(M_{M}\right)_{i j}=\frac{1}{2} \operatorname{Tr}\left(\mathbf{M}_{\mathbf{J}} \sigma_{\mathbf{j}} \overline{\mathbf{M}_{\mathbf{J}}^{\mathrm{t}}} \sigma_{\mathbf{i}}\right)
$$

## Mueller matrices and partially polarized light Time average of the previous study

## Mueller matrices are time independent

$$
\left\langle\overrightarrow{S^{\prime}}\right\rangle=\mathbf{M}_{\mathbf{M}}\langle\vec{S}\rangle
$$

$\square$
$\square$ - Polarizers along eigen axis or 1 tilted - half and quarter wave plates

## Mueller matrices and partially polarized light Time average of the previous study

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\left\langle\overrightarrow{S^{\prime}}\right\rangle=\mathbf{M}_{\mathbf{M}}\langle\vec{S}\rangle
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Mueller calculus can be extended to...

- Partially coherent light
- Cascaded optical devices


## Find the Mueller matrix of each

## Mueller matrices and partially polarized light Time average of the previous study

## Mueller matrices are time independent

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\left\langle\overrightarrow{S^{\prime}}\right\rangle=\mathbf{M}_{\mathbf{M}}\langle\vec{S}\rangle
$$

Mueller calculus can be extended to...

- Partially coherent light
- Cascaded optical devices


## Final homework

Find the Mueller matrix of each :

- Polarizers along eigen axis or $\theta$ tilted
- half and quarter wave plates
- linearly and circularly birefringent crystal

