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Introduction to Quantum Computing

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2022 - 2023



- 1 Implementation of a C-NOT gate
- 2 Example with NMR quantum computing
- 3 Molecules for quantum computing based on NMR techniques



- 1 Implementation of a C-NOT gate
 - Ising interaction
 - Two-qubit unitary evolution
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I.1. Ising interaction

Let consider an ensemble of N qubits labeled with i and j index. Those qubits interacts with an ising-type interaction (pair wise spin interaction)

$$\hat{H}_{int} = -\sum_{i,j} J_{ij} \hat{Z}_i \hat{Z}_j,$$

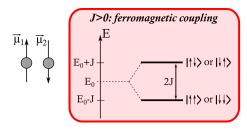
where J_{ij} is the intensity of interaction between qubits i and j.

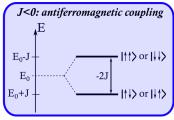
I.1. Ising interaction

It corresponds to a generic two-qubits interaction

$$\hat{H}_{int} = -J\hat{Z}_1\hat{Z}_2,$$

with J > 0 for a ferromagnetic coupling, and J < 0 for an anti-ferromagnetic coupling.





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1.2. Two-qubit unitary evolution

Applying the hamiltonian

$$\hat{H}_{int} = -J\hat{Z}_1\hat{Z}_2,$$

for a time T, the unitary evolution will be

$$\hat{C}(\gamma) = e^{-i\frac{\gamma}{2}\hat{Z}_1\hat{Z}_2},$$

with

$$\gamma = -\frac{2J}{\hbar}T.$$

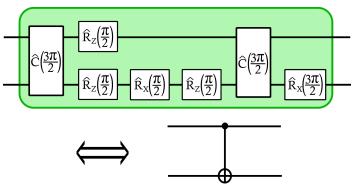
 $\hat{C}(\gamma)$ does not realize a C-NOT gate yet. Additionally, single qubit operations on each of the qubits are required to realize a C-NOT gate as follow

$$e^{-i\frac{3\pi}{4}}\hat{R}_{X2}\left(\frac{3\pi}{2}\right)\hat{C}\left(\frac{3\pi}{2}\right)\hat{R}_{Z2}\left(\frac{\pi}{2}\right)\hat{R}_{X2}\left(\frac{\pi}{2}\right)\hat{R}_{Z2}\left(\frac{\pi}{2}\right)\hat{R}_{Z1}\left(\frac{\pi}{2}\right)\hat{C}\left(\frac{3\pi}{2}\right)$$



I.2. Two-qubit unitary evolution

Circuit representation of a C-NOT gate made with a two-qubit interaction $\hat{C}\left(\gamma\right)$



Any physical two-qubit interaction that can produce entanglement can be tuned into a universal two-qubit gate (such as the C-NOT gate) when it is augmented by arbitrary single qubit operations. Ref: Bremner *et al.*, Physical Review Letter, **89**, 247902 (2002).

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II.1. NMR quantum computing

NMR is based on manipulation and measurement of nuclear spins. It's well adapted for quantum computing implementation and the first quantum calculation was implemented in NMR systems.

Reference:

Experimental realization of Shor's quantum factoring algorithm using nuclear magnetic resonance, Nature, **414**, 883-887 (2001).

 \Rightarrow factorization of 15 = 3 \times 5 in molecules by NMR.



II.1. NMR quantum computing

Introduction review articles

- NMR techniques for quantum control and computation, Rev. Mod. Phys., 76 (2004)
- Bulk spin-resonance quantum computation, Science, 275, 350 (1997)



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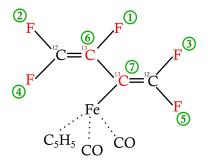
- **Qubits:** nuclear spins 1/2 in $\vec{B_0}$ ($|\uparrow\rangle$ and $|\downarrow\rangle$ as $|0\rangle$ and $|1\rangle$).
- Quantum gates: RF pulses and delay times.
- Input: Boltzmann distribution of spins (room temperature).
- Read out: detect spin states with RF coil.
- **Coherence times:** easily several seconds.

Nuclear spin in a static $\vec{B_0}$ field, with $\hat{\vec{I}}$ the nuclear spin operator and $\vec{B_0} = B_0 \vec{u_z}$,

$$\hat{\mathcal{H}}_0 = -\hbar \gamma B_0 \hat{I}_Z = -\frac{\hbar \omega_0}{2} \hat{Z},$$

where ω_0 is the Larmor frequency: precession of the qubit around \vec{u}_Z at Larmor frequency.

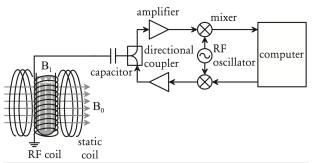
Case of $C_4F_5Fe(CH)_5(CO)_2$ - perfluorobutadienyl iron complex. Vendersypen *et al.*, Nature **414**, 883 (2001).



Nuclei labeled in red are qubits (F and ¹³C), numbered in green.

$$B_0 = 11,7T$$
.





I. L. Chuang et al., Proceedings of the Royal Society A 454, pp. 447-467 (1998).

Even without qubit/qubit coupling, the Larmor frequency depends on the atoms. Those frequencies are shifted depending on the coupling with the chemical environment

$$\mathcal{H}_0 = -\sum_{j=1}^n \hbar(1-\tilde{\sigma}_i)\gamma_i B_0 \hat{I}_Z^i = -\sum_{i=1}^n \frac{\hbar}{2} \omega_0^i (1-\tilde{\sigma}_i),$$

with ω_0^i the Larmor frequency of the nuclei i without qubit/qubit coupling and $\tilde{\sigma}_i$ the chemical frequency shift due to coupling with neighborhooding spins.

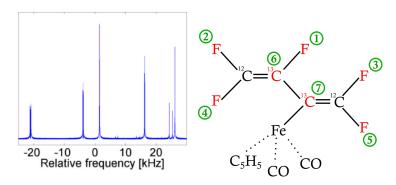
 $\tilde{\sigma}_i$ is used in analytical chemistry to determine the environment of an atom, while here it is used to address different qubits with different RF frequencies independently.

Exemple: Larmor frequency at 11.7 T for different atoms

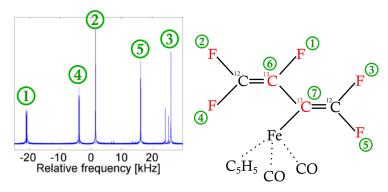
Atom	ω_0^i
¹ H	500 MHz
¹³ C	126 MHz
¹⁵ N	-51 MHz
¹⁹ F	470 MHz
³¹ P	202 MHz

with 500 MHz equivalent to 25 mK.

Chemical shifts of the five F qubits of perfluorobutandienyl iron complexe



Chemical shifts of the five F qubits of perfluorobutandienyl iron complexe



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II.3. RF field interaction: single qubit rotation

In the rotating frame and the rotating wave approximation, the hamiltonian becomes

$$\hat{H}^{rot} = -rac{\hbar\delta}{2}\hat{Z} - rac{\hbar\omega_1}{2}\left(\cos\phi\hat{X} + \sin\phi\hat{Y}
ight),$$

with $\delta=\omega_0-\omega_{RF}$ the detuning between RF and Larmor frequency, ϕ the phase of the RF field and $\omega_1=\gamma B_1$ where B_1 is the amplitude of the RF field.

For n qubits of Larmor frequencies ω_0^i , gyromagnetic factor γ_i and $\omega_1^i = \gamma_i B_1$, frequency shift $\tilde{\sigma}_i$, one notes

$$\delta_{i} = \omega_{0}^{i} \left(1 - \tilde{\sigma}_{i} \right) - \omega_{RF},$$

the detuning between RF and qubit i frequency, such that

$$\widehat{\hat{H}}^{rot} = -\sum_{i=1}^{n} \frac{\hbar}{2} \delta_{i} \hat{Z}_{i} - \frac{\hbar}{2} \sum_{i=1}^{n} \left(\cos \phi \omega_{1}^{i} \hat{X}_{i} + \sin \phi \omega_{1}^{i} \hat{Y}_{i} \right).$$



II.3. RF field interaction: single qubit rotation

Let call T the time during which the RF field is applied to the system of n qubits. It T is long enough, each spin resonance is non overlapping with the others. So if ω_{RF} is not close enough to a Larmor frequency corrected $\omega_0^i \left(1-\tilde{\sigma}_i\right)$, the effect of the RF field in negligeable. If $\omega_{RF} \approx \omega_0^i \left(1-\tilde{\sigma}_i\right)$, only the qubit i will be affected and rotate on the Bloch sphere.

If qubits have different Larmor frequencies, it is possible to manipulate each qubit individually choosing the corresponding frequency.

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When spins of different atoms are coupled, the coupling interaction energy usually results in level splitting. Consequently, the Larmor frequency of a qubit, corrected by the chemical frequency shift $\tilde{\sigma}_i$, will depend on the state of neighborhooding qubits.

In other words, the resonance RF frequency of a qubit will depends on the state of the neighborhooding qubits.

This phenomena is the key physical effect used to implement multiqubit gates.

The coupling hamiltonian might be written as follow

$$\hat{H}_J = \hbar \sum_{i < j} 4 J_{ij} \hat{I}_Z^i \hat{I}_Z^j,$$

$$\widehat{H}_J = \hbar \sum_{i < j} J_{ij} \widehat{Z}^i \widehat{Z}^j.$$

J > 0: antiferromagnetic interaction,

J < 0: ferromagnetic interaction.

Coupling constants of the five F qubits of perfluorobutandienyl iron complexe

i	$\omega_i/2\pi$	$T_{1,i}$	$T_{2,i}$	J_{7i}	J_{6i}	J_{5i}	J_{4i}	J_{3i}	J_{2i}
1	-22052.0	5.0	1.3	-221.0	37.7	6.6	-114.3	14.5	25.16
2	489.5	13.7	1.8	18.6	-3.9	2.5	79.9	3.9	
3	25088.3	3.0	2.5	1.0	-13.5	41.6	12.9		
4	-4918.7	10.0	1.7	54.1	-5.7	2.1			
5	15186.6	2.8	1.8	19.4	59.5		¹⁹ F ¹		2 19F
6	-4519.1	45.4	2.0	68.9	2		13C	13(
7	4244.3	31.6	2.0		¹⁹ F 3		6	<u> </u>	2
At A	$B_0 = 11.7 \text{ T}$:				>	¹² C ==	13C		19F
$\omega_{0,l}$	$=/2\pi = 470 \text{ N}$	ИHz					-		4
$\omega_{0,0}$	$_{\rm C}/2\pi$ = 125 M	ЛHz			¹⁹ F 5		Fe	CC_)
$[\omega_i]$	$/2\pi] = Hz, [7]$] = s, [c	/] = Hz				$C_5 H_5$	сo	

L. M. K. Vandersypen et al., Nature 414, 883 (2001)

5 resonances frequencies corresponding to 5 different qubits.

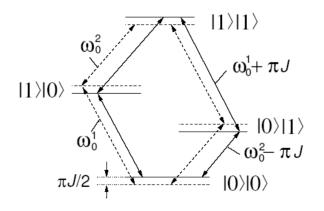


Orders of magnitude: in the case of perfluorobutadienyl iron complex at 11,7T

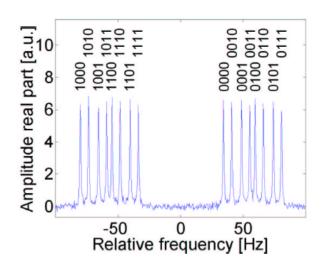
- Larmor frequency of F-type qubit \sim 470 MHz ;
- chemical frequency shift $\sim 10-20$ kHz;
- typical RF strength $\omega_1 \sim 2\pi \times 100$ kHz ;
- typical level splitting due to qubit/qubit coupling $J \sim$ few 100 Hz max.

$$\frac{100 \text{ Hz}}{470 \text{ MHz}} \sim 2 \cdot 10^{-7}, \quad \frac{100 \text{ Hz}}{10 \text{ kHz}} \sim 10^{-2}.$$











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II.5. Controlled-NOT gate in NMR

Two spins A and B as qubits. C-NOT gate flips A is B is \downarrow .

Before				
Α	В			
\uparrow	↑			
↑	\rightarrow			
\downarrow	\uparrow			
\	+			

After				
Α	В			
\uparrow				
\downarrow	\rightarrow			
	\uparrow			
↑	\rightarrow			

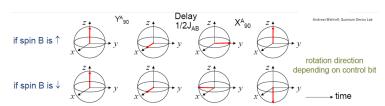
II.5. Controlled-NOT gate in NMR

Sequence:

Apply $\hat{R}_{Y,A}\left(\frac{\pi}{2}\right)$, wait for a delay

$$\Delta t = \frac{\pi}{4J_{AB}},$$

and finally apply $\hat{R}_{X,A}\left(\frac{\pi}{2}\right)$. Let describe this sequence on the Bloch sphere of spin A in the rotating frame at the larmor frequency of A, ν_A (no coupling).

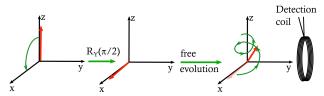


The rotation direction depends on the control qubit.

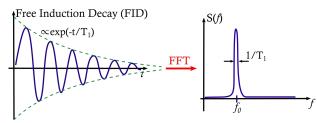
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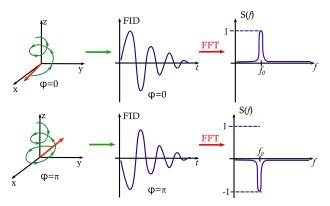
Based on NMR method. Apply $\hat{R}_Y\left(\frac{\pi}{2}\right)$ pulse and measure spin relaxation signal



A coil measures the free evolution of the spin to get access to the so-called free induction decay signal and afterwards the integrated signal or the spectrum by FFT

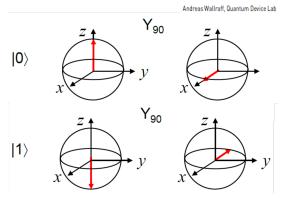


Depending on the initial position of the equator of the Bloch sphere, the phase of the free induction decay will be different



The initial angle φ on the Bloch sphere is measured by the FID signal in NMR.

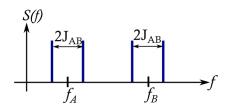
NMR is a phase sensitive detection method. A qubit state is measured with a as follow



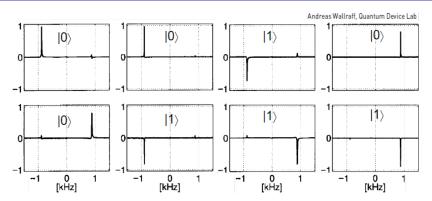
For a $|0\rangle$ state, $\varphi=0$ after the $\hat{R}_{Y}\left(\frac{\pi}{2}\right)$ pulse leading in a positive signal in the FFT of FID.

For a $|1\rangle$ state, $\varphi=\pi$ after the $\hat{R}_{Y}\left(\frac{\pi}{2}\right)$ pulse leading in a negative

Two coupled qubits in their four computational basis states have the following spectrum



NMR detection signal (FFT of FID) depends on states.



It's easier and faster experimentally to detect the sign of $S(\nu)$ rather than the frequency shift J_{AB} which is only of few hundreds of Hz.

State detection is based on sign measurement of the NMR signal $S(\nu)$.

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II.7. Example of Shor's algorithm

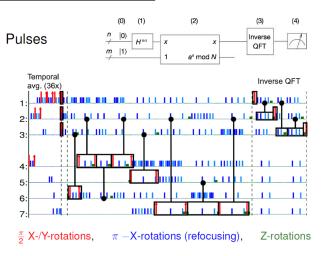
Real sequence in practice for factorization with perfluorobutadienyl iron complex.

Experimental realization of Shor's quantum factoring algorithm using nuclear magnetic resonance, Nature, **414**, 883-887 (2001).

 \Rightarrow factorization of 15 $= 3 \times 5$ in molecules by NMR with Shor's algorithm.

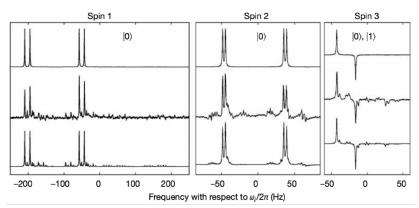
II.7. Example of Shor's algorithm

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L. M. K. Vandersypen et al., Nature 414, 883 (2001)

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II.8. Quantum state tomography

It is possible to reconstruct a state after a quantum calculation if it's not a pure state $|0\rangle$ or $|1\rangle$.

It is however possible to reconstruct a quantum state if it is possible to produce it several times. The principle consists in applying different rotations to look at the qubits from different angle.

It is called quantum state tomography.

II.8. Quantum state tomography

initial state no pulse







after X₉₀







Andreas Wallraff, Quantum Device Lab









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III.1. NMR quantum computing

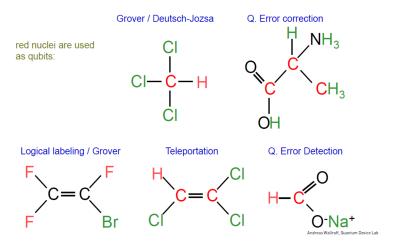
A quantum computer is a known molecule with the following desired properties

- spins 1/2 (¹H, ¹³C, ¹⁹F, ¹⁵N,...) ;
- long T_1 's and T_2 's;
- heteronuclear, or large chemical shifts (required to address different spin types independently);
- good J-coupling netword (clock-speed);
- stable, available, soluble,...

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III.2. Examples of molecules used





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